

Quantum Field Theory

Two approaches $\begin{cases} \nearrow \text{operator} \\ \searrow \text{path integral approach} \\ \text{Geometry} \end{cases}$

"Shadow" QFT : TQFT (n -dim)

1) $n-1$ -dim ^{oriented} compact manifold $\Sigma \rightarrow$
finite-dim complex vector space \mathcal{H}_Σ
"space of states"

2) n -dim. compact ^{oriented} manifold M
with boundary \rightarrow vector $Z_M \in \mathcal{H}_\Sigma$

A1 Involutivity $\mathcal{H}_{\Sigma^*} = \mathcal{H}_\Sigma^*$

A2 Multiplicativity $\mathcal{H}_{\Sigma_1 \cup \Sigma_2} = \mathcal{H}_{\Sigma_1} \otimes \mathcal{H}_{\Sigma_2}$

Remark. $\partial M = \Sigma_{\text{out}} \cup \Sigma_{\text{in}}^*$

Operator: $Z_M: \mathcal{H}_{\Sigma_{\text{in}}} \rightarrow \mathcal{H}_{\Sigma_{\text{out}}}$

Associativity: $Z_M = Z_{M_2} \circ Z_{M_1}$



A4 $\mathcal{H}_\emptyset = \mathbb{C}$ \emptyset - empty $n-1$ -manifold

Functoriality: $\phi: M \rightarrow M'$
 $\mathcal{H}_{\Sigma_{\text{in}}} \xrightarrow{Z_M} \mathcal{H}_{\Sigma_{\text{out}}}$ $\mathcal{H}_{\Sigma'_{\text{in}}} \xrightarrow{Z_{M'}} \mathcal{H}_{\Sigma'_{\text{out}}}$
 $\downarrow S(\phi_{\text{in}})$ $\downarrow S(\phi_{\text{out}})$

Unitarity: $Z_M: \mathcal{H}_{\Sigma_1} \rightarrow \mathcal{H}_{\Sigma_2}$
 $Z_{M^*}: \mathcal{H}_{\Sigma_2} \rightarrow \mathcal{H}_{\Sigma_1}$
conjugate.

A5 $Z_{\Sigma \times I} = \text{id} \mathcal{H}_\Sigma$

Corollaries

1) Take $\Sigma_2 = \Sigma_3 = \emptyset, \Sigma_2 = \Sigma$

$$Z_M = \langle Z_{M_1}, Z_{M_2} \rangle \in \mathbb{C}$$



One can take any Σ -cut:
 $M = M_1 \cup_{\Sigma} M_2$

2) Functoriality + AS

$S(\phi)$ - depends only on mapping class group of Σ

3) Construct Σ_{ϕ} by identifying boundary of cylinder $\Sigma \times [0,1]$ via $\phi \in \text{Diff}^+(\Sigma)$

$$\text{Tr } S(\phi) = Z_{\Sigma \phi}$$

In particular,
 $\dim \mathcal{H}_{\Sigma} = Z_{\Sigma \times S^1}$

Ex. 1) $n=1$. Everything is determined by a pt: $\mathcal{H}_{pt} = \mathbb{C}^N, Z_{S^1} = N$

2) $n=2$ (Dijgraaf) Frobenius algebra.

$$\mathcal{H}_{S^1} = V$$

$$(\alpha \cdot \beta, \gamma) = (\alpha, \gamma \cdot \beta)$$

- $Z \circlearrowleft$ $m: V \otimes V \rightarrow V$
- $Z \circlearrowright$ $\text{unit } 1 \in V, \iota: \mathbb{C} \rightarrow V$
- $Z \circlearrowleft$ $\text{counit } \epsilon: V \rightarrow \mathbb{C}$
- $Z \circlearrowright$ $\text{pairing } (\cdot, \cdot): V \otimes V \rightarrow \mathbb{C}$
- $Z \circlearrowright$ $\text{coproduct } \Delta: V \rightarrow V \otimes V$
- $Z \text{ closed surface of genus } g$ $= \epsilon \circ (\text{mod } \Delta)^g \circ 1$

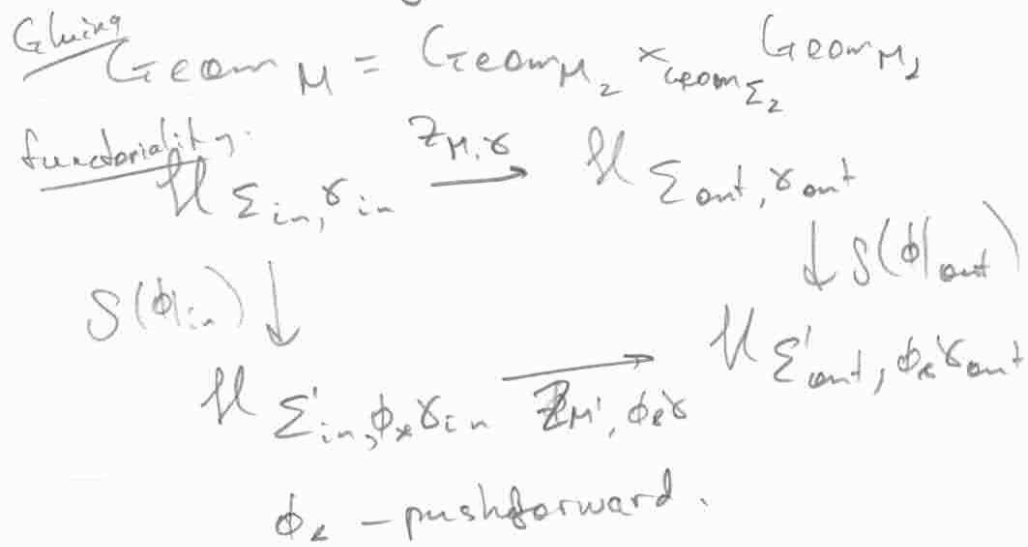
3) $n \geq 3$ Hard. Needed a refinement

Ex. Cohomology, quantum cohomology of comp. X
 $\mathbb{P}^n, X^{n+2} = 0 \rightarrow q, \alpha \cdot \beta = \alpha \wedge \beta, \int \alpha \wedge \beta = \eta(\alpha, \beta)$

Segal's QFT

Manifold with some geometric data: Riemann metric, complex structure, etc

Modification of Atiyah axioms to include geom. data:



Ex. Quantum mechanics

\pm orientation for pts:

$+$ $\rightarrow \mathcal{H}$
 $-$ $\rightarrow \mathcal{H}^*$

Metric on d -dim. cobordisms

$Z(t) = Z_{\mathbb{I}^+} \in \text{End}(\mathcal{H})$

interval of length t

$Z(t_1 + t_2) = Z(t_1)Z(t_2)$

Assume theory is unitary: 3
Unitarity: $Z_M: \mathcal{H}_{\Sigma_1} \rightarrow \mathcal{H}_{\Sigma_2}$ } conjugate
 $Z_{M^*}: \mathcal{H}_{\Sigma_2} \rightarrow \mathcal{H}_{\Sigma_1}$

$Z_M^*(t) = Z_M(t)$ for cylinder
 $Z(t) = e^{-\frac{i}{\hbar} H t}$ (Hamiltonian)

$\mathcal{H} = \mathcal{L}^2(X)$ $\psi(x) \rightarrow \int_{x \in Y} dy \psi(t, x, y) \psi^\dagger(y)$
↑
operator.

Path integrals

Classical field theory on n -manifold

i) Space of fields:

$F_M = \Gamma(M, \mathcal{F}_M)$ - sections of some "sheaf"

1) functions

2) sections of vector bundles.

3) Maps: $M \rightarrow N$ target space
 "sigma-models"

4) Connections on principal G -bundles over M
 Chern-Simons Yang-Mills

ii) Action functional $S_M: F_M \rightarrow \mathbb{R}$
 $S_M(\phi) = \int_M \mathcal{L}(\phi)$ - Lagrangian density

Ex. Quantum mechanics in \mathbb{R}^N

$$S_M(\gamma) = \int (\sum_i \frac{\dot{q}_i^2}{2} - V(q)) dt$$

$\gamma: \mathbb{R} \rightarrow \mathbb{R}^N$ \downarrow \hat{V} potential
 $t \rightarrow q_i(t)$

Path integral quantization:

$$Z_M(\hbar) = \int_{F_M} D\phi e^{\frac{i}{\hbar} S_M(\phi)}$$

Hard to define. Easy to consider asymptotic expansion $\hbar \rightarrow 0$.

Take M with boundary B_Σ -boundary values of fields.

Take pullback $\Sigma \xrightarrow{i} M$ $\phi \xrightarrow{i^*} \phi|_\Sigma$

$$\mathcal{H}_\Sigma = \text{Func}(B_\Sigma)$$

$$Z_M(\phi_\Sigma, \hbar) = \int_{\phi \in F_M, \phi|_{\partial M} = \phi_\Sigma} D\phi e^{\frac{i}{\hbar} S_M(\phi)}$$

Argument for gluing:

$\mathbb{R}^1/2$

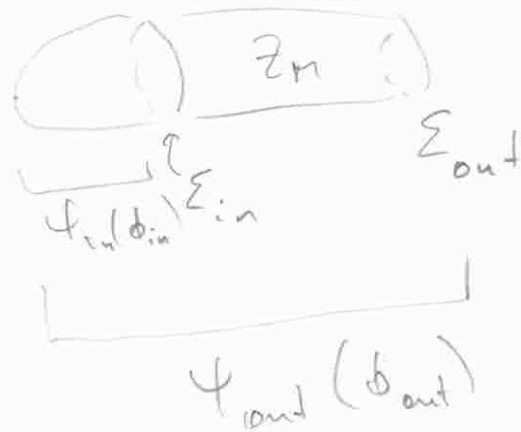
$$Z_M = \int D\phi_\Sigma Z_{M_1}(\phi_\Sigma) Z_{M_2}(\phi_\Sigma)$$

$$B_\Sigma \supset \phi_\Sigma$$



$$M = M_1 \cup_\Sigma M_2$$

$$\langle \psi_{\text{out}} | \psi_{\text{in}} \rangle = \int D\phi_{\text{in}} Z_M(\phi_{\text{out}}, \phi_{\text{in}}) \psi_{\text{in}}(\phi_{\text{in}})$$



Path integral - making sense

$$f \in C^\infty(x) \quad \int_x \mu e^{\frac{i}{\hbar} f(x)}$$

Stationary phase formula:

$$\int_x \mu e^{\frac{i}{\hbar} f(x)} \underset{\hbar \rightarrow 0}{\sim} \sum_{x_0 \in \text{crit. points of } f} e^{\frac{i}{\hbar} f(x_0)} |\det f''(x_0)|^{-1/2} e^{\frac{\pi i}{4} \sum \text{sign } f''(x_0)} \left(\frac{\dim x}{2}\right) + o(\hbar^{n/2})$$

positive - negative eigens.

$\Delta f|_{x_0} = 0$ - 'stationary phase points'
oscillations slow down

Improvement:

$$\int_x \mu e^{\frac{i}{\hbar} f(x)} \underset{\hbar \rightarrow 0}{\sim} \sum_{x \in \text{crit. pt. of } f} e^{\frac{i}{\hbar} f(x_0)} |\det f''(x_0)|^{-1/2} e^{\frac{\pi i}{4} \sum \text{sign } f''(x_0)} \left(\frac{\dim x}{2}\right)$$

vertices of valency 3

$\chi(\Gamma) \leq 0$
euler char. of graph.

$$\sum_{\Gamma} \frac{\hbar^{-\chi(\Gamma)}}{|\text{Aut}(\Gamma)|} \Phi_{\Gamma}$$

Feynman diagrams

half-edges - decorated w indices

i) Every edge:
 $\rightarrow f''(x_0)_{i_1, i_2}$ - inverse of Hessian

ii) vertex:
 $\partial_{i_1} \dots \partial_{i_k} f(x_0)$ - k-th partial derivative at crit. pt.

$$\Phi_{\Gamma} = i^{E+V} \sum_{i_1, \dots, i_{2E}} \prod_{\text{edges } e=(h_i, h_j)} f''(x_0)_{i_1, i_2} \prod_{\text{vertices } v} \partial_{i_1} \dots \partial_{i_{\text{val}(v)}} f(x_0)$$

Ex.



Θ -graph

$\chi(\Theta) = -1$

$\mathbb{Z}_2 \times S_2$
symmetry group

$\frac{\hbar}{12}$ - coeff.

$$\Phi \left(\begin{array}{ccc} i & & \\ | & & \\ j & - & m \\ | & & \\ k & & n \end{array} \right) = i^{3+2} \times$$

$$\sum_{i, j, k, l, m, n} f''(x_0)_{i, j}^{-1} f''(x_0)_{j, m}^{-1} f''(x_0)_{m, n}^{-1} f''(x_0)_{k, l}^{-1} f''(x_0)_{l, i}^{-1} f''(x_0)_{n, k}^{-1}$$

Problem: "Gauge symmetry"
 Instead of isolated fixed pts
 There is a continuous symmetry
 e.g. gauge group G . How to
 isolate G -orbits?

BV-formalism:

$\{F, S\} \rightarrow$ new data:

i) \mathcal{F} graded supermanifold \mathcal{F}
 "space of BV-fields"
 ω - odd symplectic structure of $\text{deg} = 1$

ii) S_{BV} on \mathcal{F} :

$\{S_{BV}, S_{BV}\} = 0$ (deg 1 Poisson bracket)
 $\{S_{BV}, -\} = Q$
 $Q^2 = 0 \Rightarrow \mathcal{F}$ -dgmanifold

$\int_{\mathcal{F}} e^{\frac{i}{\hbar} S_{BV}} \rightarrow \int_{\mathcal{L} \subset \mathcal{F}} e^{\frac{i}{\hbar} S_{BV}}$
 Lagrangian submanifold.

Relative TQFT:

Example: M - 3d manifold
 $\Delta \subset M$ - oriented 1-d im submanifold
 $\partial M = \Sigma$ assume Δ transversal to Σ
 and $+$ - pts

(G, \mathcal{F}) - Indices
 Jones-Witten CS repres of G